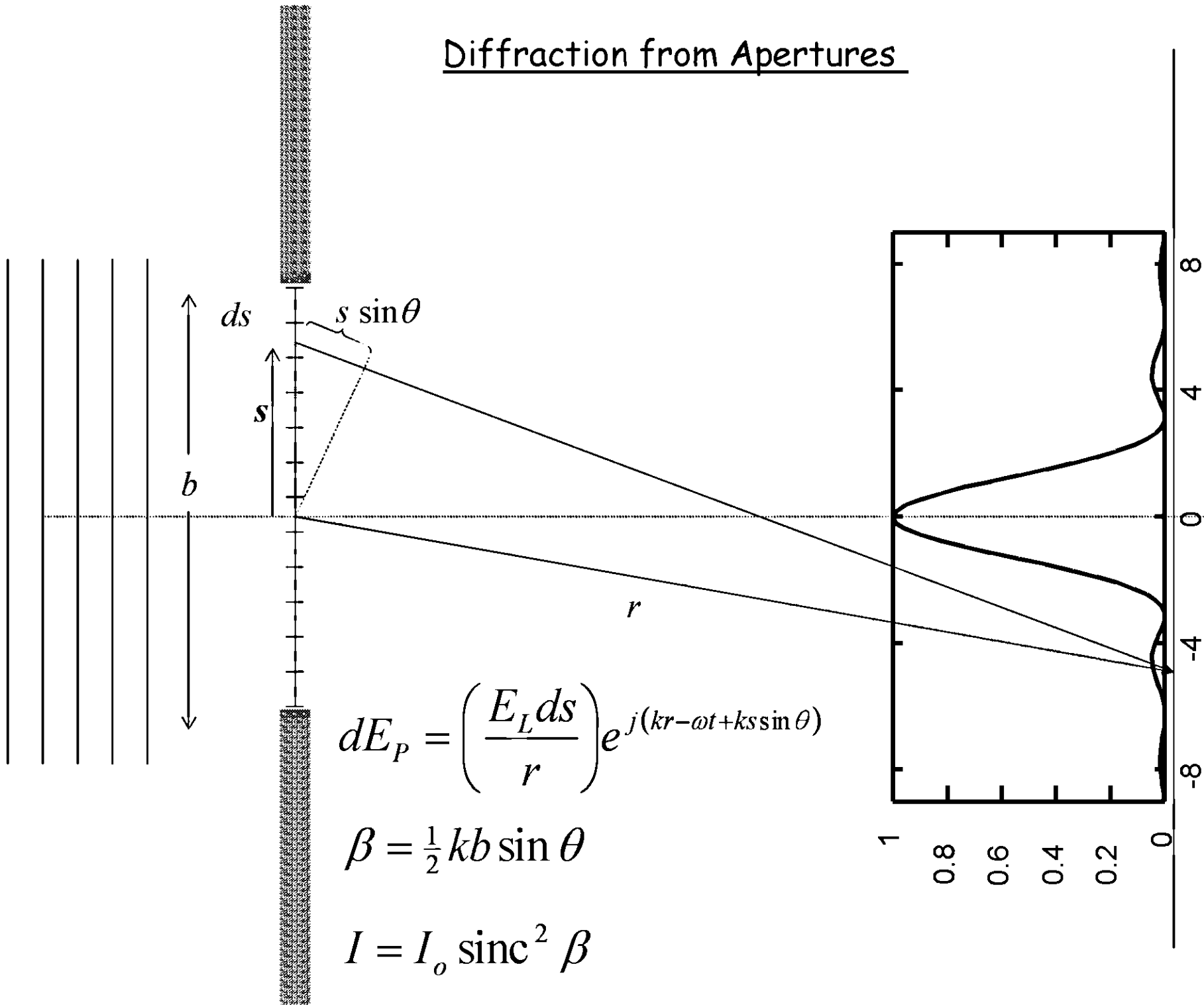
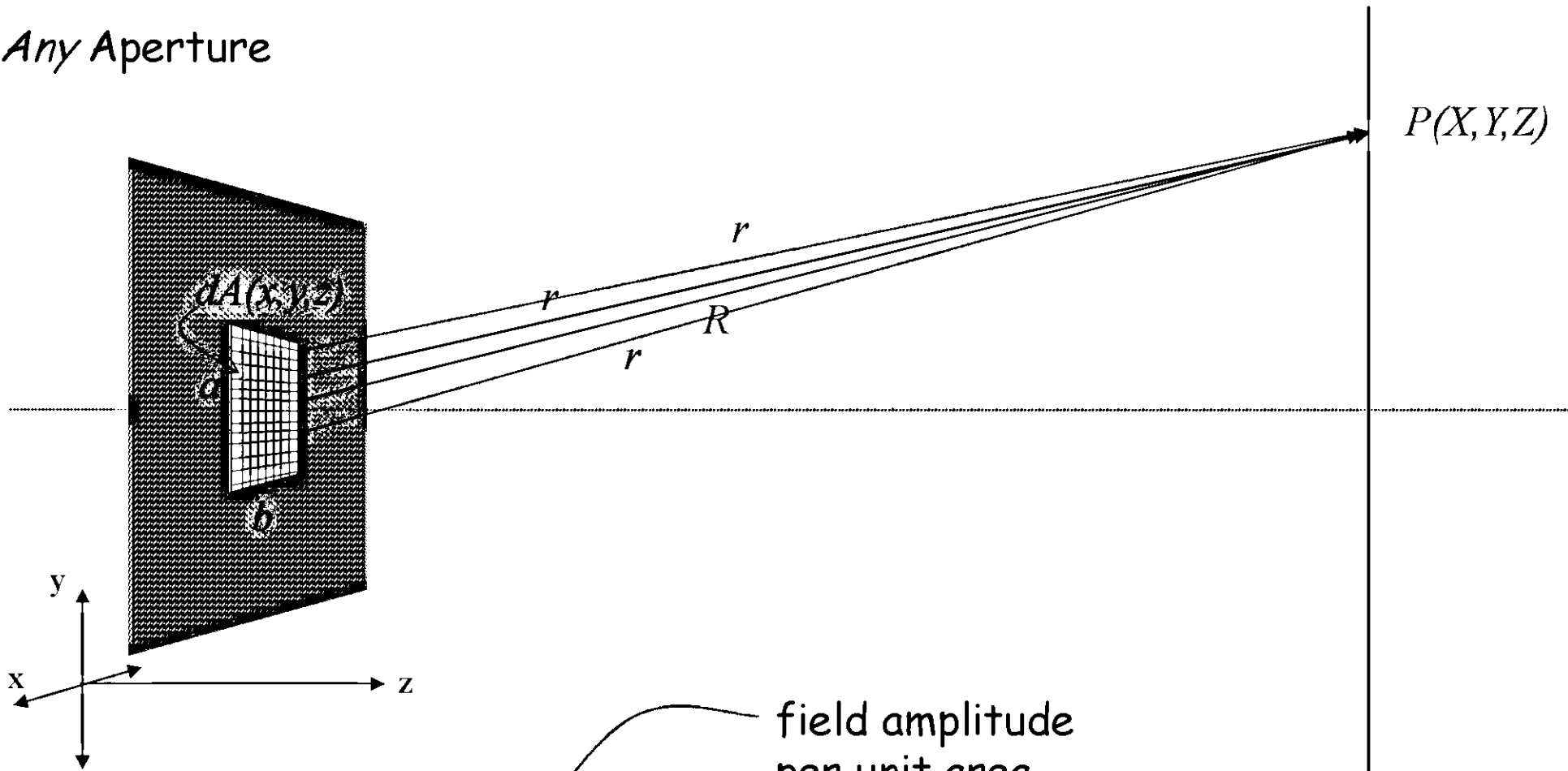


Diffraction from Apertures



Any Aperture



field amplitude
per unit area

$$dE_P = \left(\frac{E_A dA}{r} \right) e^{j(\omega t - kr)}$$

Phase differences
in r 's

$$R = \sqrt{X^2 + Y^2 + Z^2}$$

$$r = \sqrt{(X - x)^2 + (Y - y)^2 + Z^2}$$

$$r \approx R$$

They are close enough for the amplitude term
(the difference is small compared to the screen distance)

$$dE_p = \left(\frac{E_A dA}{R} \right) e^{j(\omega t - kr)}$$

They are *not* close enough for the phase term
(the difference is *not* small compared to the wavelength)

$$r = \sqrt{X^2 + Y^2 + Z^2 - 2Xx + x^2 - 2Yy + y^2}$$

$$r = R\sqrt{1 - 2(Xx + Yy)/R^2 + (x^2 + y^2)/R^2}$$

$$r \approx R\sqrt{1 - 2(Xx + Yy)/R^2}$$

binomial expansion:

$$r \approx R\left(1 - (Xx + Yy)/R^2\right)$$

$$dE_p = \left(\frac{E_A dA}{R}\right) e^{j(\omega t - kR[1 - (Xx + Yy)/R^2])}$$

$$dE_p = \left(\frac{E_A dA}{R} \right) e^{j(\omega t - kR [1 - (Xx + Yy)/R^2])}$$

$$dE_p = \left(\frac{E_A dA}{R} \right) e^{j(\omega t - kR)} e^{jk(Xx + Yy)/R}$$

$$E_p = \left(\frac{E_A}{R} \right) e^{j(\omega t - kR)} \iint_{\text{aperture}} e^{jk(Xx + Yy)/R} dA \quad \text{For any shape!}$$

Rectangular:

$$E_p = \left(\frac{E_A}{R} \right) e^{j(\omega t - kR)} \int_{-\frac{a}{2}}^{\frac{a}{2}} \int_{-\frac{b}{2}}^{\frac{b}{2}} e^{jk(Xx + Yy)/R} dx dy$$

$$E_P = \left(\frac{E_A}{R} \right) e^{j(\omega t - kR)} \int_{-\frac{a}{2}}^{\frac{a}{2}} e^{\frac{jkYy}{R}} dy \int_{-\frac{b}{2}}^{\frac{b}{2}} e^{\frac{jkXx}{R}} dx$$

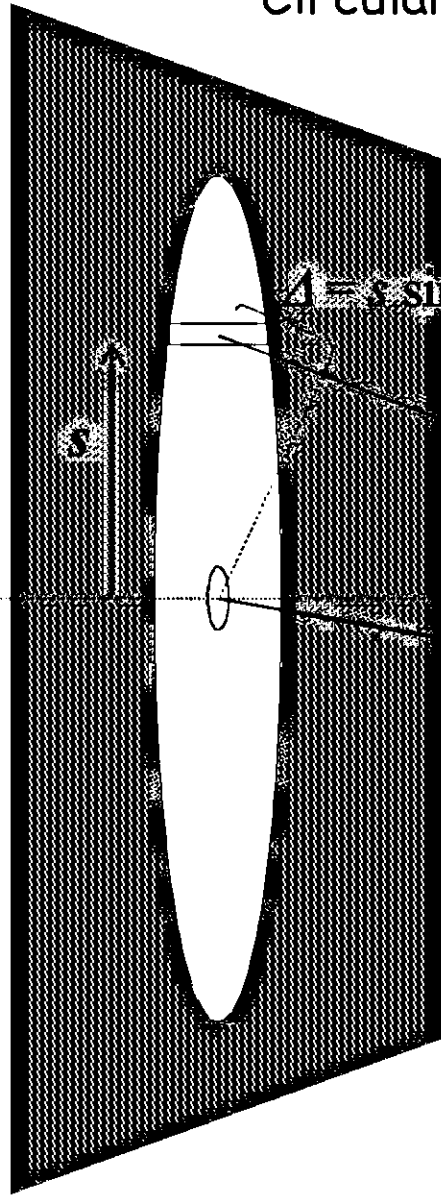
$$\beta = \frac{1}{2} kb \sin \theta = \frac{1}{2} kb \frac{Y}{R} \qquad \alpha = \frac{1}{2} ka \sin \theta = \frac{1}{2} ka \frac{X}{R}$$

From here its just like the single slit:

$$I(X, Y) = I_o \text{sinc}^2 \beta \text{sinc}^2 \alpha$$

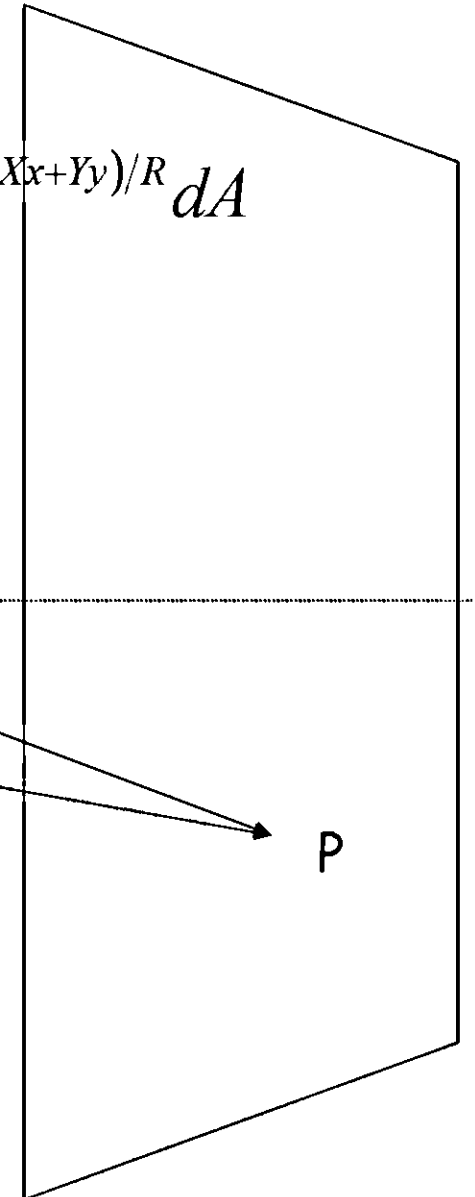
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Circular Apertures

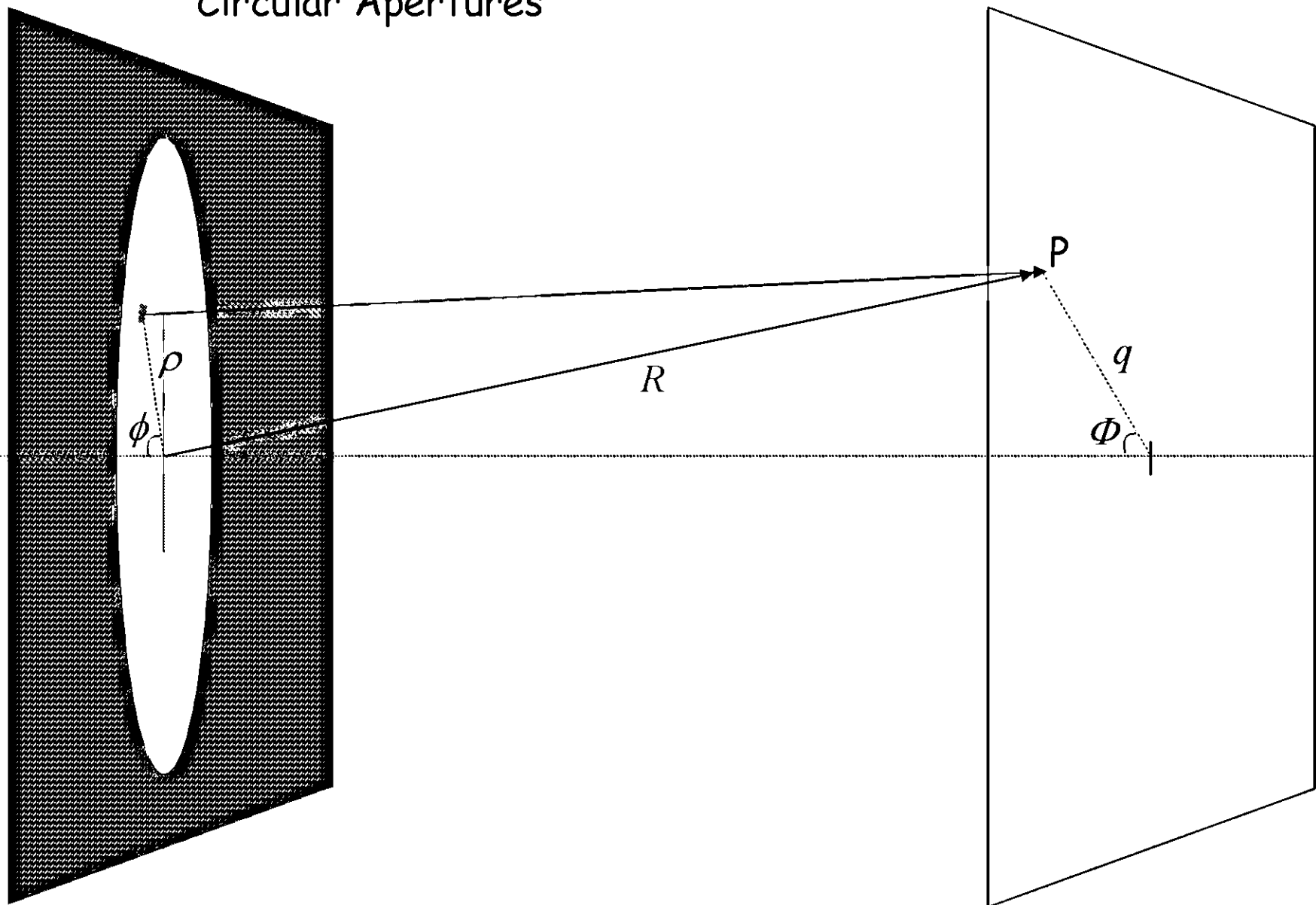


$$E_P = \left(\frac{E_A}{R} \right) e^{j(\omega t - kR)} \iint_{\text{aperture}} e^{jk(Xx+Yy)/R} dA$$

$$E_P = \frac{E_A}{r} e^{j(kr - \omega t)} \int_A e^{j(ks \sin \theta)} dA$$



Circular Apertures



$$E_P = \left(\frac{E_A}{R} \right) e^{j(\omega t - kR)} \int_{\rho=0}^a \int_{\phi=0}^{2\pi} e^{jk(\rho q/R)\cos(\phi-\Phi)} \rho d\phi d\rho$$

Φ is arbitrary, might as well make it zero:

$$E_P = \left(\frac{E_A}{R} \right) e^{j(\omega t - kR)} \int_{\rho=0}^a \int_{\phi=0}^{2\pi} e^{jk(\rho q/R)\cos(\phi)} d\phi \rho d\rho$$

$$x = k\rho q/R$$

$$E_P = \left(\frac{E_A}{R} \right) e^{j(\omega t - kR)} \int_{\rho=0}^a \boxed{\int_{\phi=0}^{2\pi} e^{jx\cos(\phi)} d\phi} \rho d\rho$$

Bessel Function!

Bessel function of the first kind (J) of order zero (0)

$$J_0(x) = \frac{1}{2\pi} \int_0^{2\pi} e^{jx \cos \phi} d\phi$$

Bessel functions $J_m(x)$ solve this differential equation:

$$x^2 \frac{d^2 y}{dx^2} + x \frac{dy}{dx} + (x^2 - m^2)y = 0$$

Recurrence relations let you solve integrals and derivatives:

$$\frac{d}{dx} [x J_1(x)] = x J_0(x) \quad x J_1(x) = \int_0^x x' J_0(x') dx'$$

They can be expressed as series:

$$J_0(x) = \sum_{n=0}^{\infty} \frac{-1^n}{n! \Gamma(n+1)} \left(\frac{1}{2} x\right)^{2n}$$

$$E_p = 2\pi \left(\frac{E_A}{R} \right) e^{j(\omega t - kR)} \int_{\rho=0}^a J_0(x) \rho d\rho$$

$$E_p = 2\pi \left(\frac{E_A}{R} \right) e^{j(\omega t - kR)} \int_{x=0}^{x=kaq/R} J_0(x) \frac{xR}{kq} \frac{R}{kq} dx$$

$$E_p = 2\pi \left(\frac{E_A}{R} \right) e^{j(\omega t - kR)} \left(\frac{R}{kq} \right)^2 (kaq/R) J_1(kaq/R)$$

$$E_p = 2 \left(\frac{E_A A}{R} \right) e^{j(\omega t - kR)} \left(\frac{R}{kaq} \right) J_1(kaq/R)$$

$$I_P = \frac{1}{2} \epsilon_0 c^4 \left(\frac{E_A A}{R} \right)^2 \left[\frac{J_1(kaq/R)}{kaq/R} \right]^2$$

$$I_P = I_o \left[\frac{2J_1(kaq/R)}{kaq/R} \right]^2$$

